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Papaefstathiou, Andreas ; Sakurai, Kazuki ; Takeuchi, Michihisa

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Higgs boson to di-tau channel in chargino-neutralino searches at the LHC

Andreas Papaefstathiou,^a Kazuki Sakurai^b and Michihisa Takeuchi^b

^a*Physik Institut, Universität Zürich,
Switzerland*

^b*Department of Physics, Theoretical Particle Physics & Cosmology,
King's College London,
United Kingdom.*

E-mail: andreasp@physik.uzh.ch, kazuki.sakurai@kcl.ac.uk,
michihisa.takeuchi@kcl.ac.uk

ABSTRACT: We consider chargino-neutralino production, $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow (h \tilde{\chi}_1^0)(W^\pm \tilde{\chi}_1^0)$, which results in Higgs boson final states that subsequently decay (inclusively) to leptons (either $h \rightarrow \tau^+ \tau^-$ or $h \rightarrow W^+ W^- \rightarrow (e^+ e^-, \mu^+ \mu^-, \tau^+ \tau^-) + \cancel{E}_T$). Such channels are dominant in large regions of the allowed supersymmetric parameter space for many concrete supersymmetric models. The existence of leptons allows for good control over the backgrounds, rendering this channel competitive to the conventional $h \rightarrow b\bar{b}$ channel. We include hadronic decays of the τ leptons in our analysis through a τ -identification algorithm. We consider integrated luminosities of 100 fb^{-1} , 300 fb^{-1} and 3000 fb^{-1} , for an LHC running at pp centre-of-mass energy of 14 TeV and provide the expected constraints on the M_2 – M_1 plane.

KEYWORDS: Supersymmetry Phenomenology, Hadronic Colliders

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1 Introduction

One of the primary goals of the CERN Large Hadron Collider (LHC) is to discover or rule out weak scale supersymmetry (SUSY). So far the ATLAS and CMS collaborations have conducted a number of direct SUSY searches in many different channels. The absence of excesses in those channels over the Standard Model (SM) background in turn placed impressive constraints on the SUSY parameter space. The limit is particularly stringent for coloured SUSY particles because of their large production cross sections. For instance, gluino and light flavour squarks are excluded up to masses of about

1 – 1.5 TeV [1–7], although the precise mass bounds depend on the details of the decay chains and mass spectrum.¹

The recent observation of a SM-like Higgs boson [12, 13] also provides interesting implications and opportunities for the exploration of SUSY phenomenology. First of all, the observed mass ~ 125 GeV and the measured properties of the SM-like Higgs boson are consistent with the lightest CP-even Higgs (h) in the minimal SUSY extension of the SM (MSSM) especially when the masses of scalar superparticles are larger than the several TeV [14–17]. Such scenarios are also consistent with the null results of direct SUSY searches and the precise measurements of flavour-changing neutral currents (FCNC) and CP-violating observables.

Even though the scalars are anticipated to be heavy, it is possible to have relatively light gauginos in the SUSY spectrum. In particular, the electroweak (EW) gauginos can exist and still be very light, since their production cross sections are much smaller than the coloured SUSY particles of the same mass. Indeed, in concrete models, the EW gauginos tend to be much lighter than the coloured SUSY particles. This is due to the fact that the renormalisation group evolution (RGE) increases coloured SUSY particle masses at low energies, owing to their strong QCD interaction, whilst the effect is much smaller for EW gauginos. It is known [18] that if the gaugino GUT relation ($M_3 : M_2 : M_1 \sim 7 : 2 : 1$) holds, the production of EW gauginos can dominate over gluino pair production at the 14 TeV LHC due to the mass hierarchy. Moreover, many SUSY breaking scenarios predict a large mass splitting between gauginos and scalars [19–24]. Unlike the scalar masses, gaugino mass terms are prohibited by R-symmetry, and their mass generation mechanism may be very different. In the scenarios where R-symmetry is only weakly broken, the gauginos tend to be much lighter than the scalars. In such scenarios, gauginos are the only SUSY particles which are accessible at the LHC [25–28].

The EW gauginos, namely, charginos and neutralinos, have already been intensively searched for at the LHC. ATLAS and CMS interpreted their results in the context of simplified models, where several assumptions were made. For instance, the lightest neutralino ($\tilde{\chi}_1^0$) was assumed to be bino-like and the second lightest neutralino ($\tilde{\chi}_2^0$) and the lighter chargino ($\tilde{\chi}_1^\pm$) wino-like, while possessing the same mass, $m_{\tilde{\chi}_2^0} = m_{\tilde{\chi}_1^\pm}$. In these simplified models, particular decays of the chargino and the second lightest neutralino with 100% branching ratios were considered. The most stringent constraints were found for the models where the $\tilde{\chi}_2^0$ and $\tilde{\chi}_1^\pm$ decay exclusively into on-shell sleptons ($\tilde{\ell}$ and $\tilde{\nu}$). In this case, $m_{\tilde{\chi}_2^0} = m_{\tilde{\chi}_1^\pm}$ is excluded up to about 700 GeV with $m_{\tilde{\chi}_1^0} \lesssim 300$ GeV [29, 30]. Simplified models with the chargino and neutralino decays leading to di- τ final states via on-shell $\tilde{\tau}$ and $\tilde{\nu}_\tau$ have also been searched for, and the limit was found to be $m_{\tilde{\chi}_2^0} = m_{\tilde{\chi}_1^\pm} \gtrsim 300$ (350) GeV with $m_{\tilde{\chi}_1^0} \lesssim 100$ (50) GeV [29, 30]. If the sleptons and staus are heavier than the EW gauginos, the $\tilde{\chi}_1^\pm$ predominantly decays to W^\pm and $\tilde{\chi}_1^0$. On the other hand, the $\tilde{\chi}_2^0$ has two possible decay modes: $\tilde{\chi}_2^0 \rightarrow Z\tilde{\chi}_1^0$ and $\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0$. The former has been searched for and the resulting limit was $m_{\tilde{\chi}_2^0} = m_{\tilde{\chi}_1^\pm} \gtrsim 350$ with $m_{\tilde{\chi}_1^0} \lesssim 100$ GeV [29]. The latter process, $\tilde{\chi}_2^0\tilde{\chi}_1^\pm \rightarrow (h\tilde{\chi}_1^0)(W^\pm\tilde{\chi}_1^0)$, has also been looked for recently by employing the $h \rightarrow b\bar{b}$ channel.

¹See, e.g., [8–11] for the recent ideas and programmes to address this problem.

lower bound on $m_{\tilde{\chi}_2^0} = m_{\tilde{\chi}_1^\pm}$	search hypothesis
700 GeV [29, 30]	$\tilde{\chi}_2^0, \tilde{\chi}_1^\pm$ both decay exclusively to sleptons, $m_{\tilde{\chi}_1^0} = 300$ GeV.
300 (350) GeV [29, 30]	$m_{\tilde{\chi}_1^0} \lesssim 100$ (50) GeV.
350 GeV [29]	staus heavier than EW gauginos, $\tilde{\chi}_1^\pm \rightarrow W^\pm \tilde{\chi}_2^0, \tilde{\chi}_2^0 \rightarrow Z \tilde{\chi}_1^0, m_{\tilde{\chi}_1^0} \lesssim 100$ GeV.
200 GeV [31], 300 GeV [32]	staus heavier than EW gauginos, $\tilde{\chi}_1^\pm \rightarrow W^\pm \tilde{\chi}_2^0, \tilde{\chi}_1^\pm \rightarrow h \tilde{\chi}_1^0$, with $h \rightarrow b\bar{b}$, when $m_{\tilde{\chi}_1^0} \lesssim 30$ GeV.

Table 1. Summary of the EW gaugino searches as given in the main text.

This channel suffers from an overwhelmingly large $t\bar{t}$ background and only weak constraints have been found. The bound is $m_{\tilde{\chi}_2^0} = m_{\tilde{\chi}_1^\pm} \gtrsim 200$ GeV [31] and 300 GeV [32] only when $m_{\tilde{\chi}_1^0} \lesssim 30$ GeV. A summary of EW gaugino searches is given in Table 1 for convenience.

The fact that current searches provide weak constraints is not the only reason the $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow (h \tilde{\chi}_1^0)(W^\pm \tilde{\chi}_1^0)$ process is especially interesting for further study. Firstly, in this process one can take advantage of the discovery of the SM-like Higgs boson, making use of the of its properties as measured in the present dataset [33]. Identifying the observed boson as the lightest CP-even Higgs in the MSSM allows us to make a precise prediction of the $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow (h \tilde{\chi}_1^0)(W^\pm \tilde{\chi}_1^0)$ signature, which is necessary for the limit calculation and also useful in designing optimal search strategies for this mode. Secondly, as we will see in section 2, the scenarios with heavy scalar SUSY particles may imply that $\tilde{\chi}_2^0$ predominantly decays into h and $\tilde{\chi}_1^0$.

In this paper, we study the exclusion and discovery reach of the $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow (h \tilde{\chi}_1^0)(W^\pm \tilde{\chi}_1^0)$ process, using the W decays to electrons muons or taus and the $h \rightarrow \tau\tau$ and $h \rightarrow WW \rightarrow (\tau/\ell, \nu)(\tau/\ell, \nu)$ modes. Our study differs from earlier studies for $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow (h \tilde{\chi}_1^0)(W^\pm \tilde{\chi}_1^0)$ [18, 33–35], which have focused on the decays of the W to electrons or muons and $h \rightarrow b\bar{b}$ modes. The obvious advantage of the channel with $h \rightarrow b\bar{b}$ is its relatively large branching ratio $BR(h \rightarrow b\bar{b})$. However, this channel suffers from an overwhelming $t\bar{t}$ background. Employing $h \rightarrow \tau\tau$ and $h \rightarrow WW \rightarrow (\tau/\ell, \nu)(\tau/\ell, \nu)$ introduces a reduction of the branching ratio, by a factor of $[BR(h \rightarrow \tau\bar{\tau}) + BR(h \rightarrow WW \rightarrow (\tau/\ell, \nu)(\tau/\ell, \nu))]/BR(h \rightarrow b\bar{b}) \sim 0.15$, but the $t\bar{t}$ background can be reduced significantly by vetoing b -jets and requiring two τ s in the final state as we will see in section 3. We will demonstrate that the channel with $h \rightarrow \tau\tau$ and $h \rightarrow WW \rightarrow (\tau/\ell, \nu)(\tau/\ell, \nu)$ can provide competitive discovery and exclusion prospects to those obtained in the channel with the $h \rightarrow b\bar{b}$ mode.

Revealing the details of the EW gaugino sector is especially important. It is commonly believed that this sector contains the particle that can be a candidate for dark matter. Moreover, studying the accessible mass scale of the EW gauginos at the LHC is important [36] for the planning of future collider programmes.

The article is organised as follows: in the next section, we provide the details of the setup we use for the $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow (h\tilde{\chi}_1^0)(W^\pm \tilde{\chi}_1^0)$ mode and discuss the cross section and branching ratio of EW gauginos with particular attention to heavy scalar scenarios with a large μ -term. In section 3 we provide details of the Monte Carlo simulation performed to generate the samples used in the analysis and give details of the algorithm employed for the identification of jets originating from hadronic decays of τ leptons. We then proceed to outline the details of our discrimination analysis, which forms the basis for defining the signal regions for the SUSY parameter space scan. The results of the parameter space scan are presented and discussed in section 4. We conclude in section 5. Supplementary appendices describe the definition of a kinematic variable used in our analysis, the calculation of cross sections for signal and background and statistical methods with low event numbers. The last appendix in particular describes a systematic way to recast our results onto the other scenarios. The application includes higgsino NSLP scenarios with a bino LSP and higgsino/wino NLSP scenarios with a gravitino LSP as discussed, for example, in [37–40].

2 The $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow (h\tilde{\chi}_1^0)(W^\pm \tilde{\chi}_1^0)$ mode

In this section we describe the setup of our analysis and clarify the assumptions we made in the chargino and neutralino sectors. Moreover, we discuss the cross sections and branching ratios of the production and decay modes relevant to our analysis.

2.1 The setup

Throughout this paper we consider CP-conserving EW gaugino sector and assume $m_{\tilde{\chi}_2^0} \simeq m_{\tilde{\chi}_1^\pm} > m_{\tilde{\chi}_1^0}$ for simplicity. This relation is realised in many SUSY breaking scenarios, particularly in the cases where $|\mu| \gg M_2 > M_1$ and $M_2 \gg |\mu| > M_1$. The former case is motivated by the heavy scalar scenario. In the MSSM, the soft scalar masses for H_u and H_d and the μ -parameter are related by the EW symmetry breaking condition [41]

$$\frac{m_Z^2}{2} = \frac{m_{H_d}^2 - m_{H_u}^2 \tan^2 \beta}{\tan^2 \beta - 1} - |\mu|^2. \quad (2.1)$$

In the heavy scalar scenario, the scale of m_{H_u} and m_{H_d} is very large ($m_{H_{u/d}} \gg M_2$). In this case, the condition (2.1) implies that the μ -parameter is expected to be of the same scale as the scalar masses, unless m_{H_u} and m_{H_d} are carefully tuned at the EW scale in such a way that the first terms in the right hand side of eq. (2.1) becomes unnaturally small.²

In this section we assume the scale of μ is equal to the scalar masses and $|\mu| \gg M_2 > M_1 > 0$. However, the collider analysis described in section 3 is applicable to other scenarios as far as the $\tilde{N}\tilde{C}^\pm \rightarrow (h\chi)(W^\pm \chi)$ topology is concerned, where \tilde{N} and \tilde{C}^\pm are massive BSM particles with the same mass and χ is an invisible particle with an arbitrary mass. One such scenario involves a bino LSP scenario with a higgsino NLSP, $M_2 \gg |\mu| > M_1$. The application also includes gravitino LSP scenarios with wino or higgsino NSLP as discussed for example in [37–40], where the same topology is realised by $\tilde{\chi}_1^0 \tilde{\chi}_1^\pm \rightarrow (h\tilde{G})(W^\pm \tilde{G})$ with \tilde{G} being gravitino. We will get back to this point in the end of this section.

²Even in that case, the same size of tuning is required on the μ -parameter.

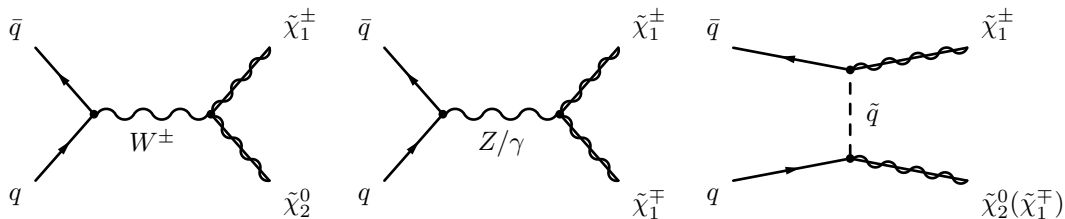


Figure 1. The tree-level diagrams for the relevant $\tilde{\chi}_2^0$ and $\tilde{\chi}_1^\pm$ production.

2.2 The cross sections

Figure 1 shows the tree-level diagrams for the relevant modes of $\tilde{\chi}_2^0$ and $\tilde{\chi}_1^\pm$ production. There are two types of diagrams which may interfere: s -channel diagrams with gauge boson exchange and t -channel diagrams with squark exchange. The t -channel diagrams are suppressed by the squark mass and it is expected that the contribution of this diagram decreases as the squark mass increases.

Figure 2 shows the NLO production cross sections for the $\tilde{\chi}_2^0\tilde{\chi}_1^\pm$ and $\tilde{\chi}_1^+\tilde{\chi}_1^-$ modes at the 14 TeV LHC as functions of the squark mass. The cross sections have been calculated using Prospino 2.1 [42, 43] with all the charges summed. In the plot and throughout the paper, we take $|\mu| = m_{\tilde{q}}$ for simplicity. For the specific plot, we take $M_2 = 350$ GeV and $M_1 = 100$ GeV. The solid and dashed curves correspond to $\tan\beta = 2$ and 50, respectively. As a result of destructive interference between the s -channel gauge boson exchange diagram and the t -channel squark exchange diagram, the $\tilde{\chi}_2^0\tilde{\chi}_1^\pm$ and $\tilde{\chi}_1^+\tilde{\chi}_1^-$ production cross sections increase as the squark mass increases. For a squark mass larger than ~ 4 TeV, the contribution of the squark exchange diagram is decoupled and the cross sections become insensitive to the squark mass. It is interesting to note that the $\tilde{\chi}_2^0\tilde{\chi}_1^\pm$ and $\tilde{\chi}_1^+\tilde{\chi}_1^-$ cross sections are maximised in the limit of large squark mass. This gives additional motivation to perform EW gaugino searches in the context of heavy scalar scenarios.

Figure 3 shows the NLO cross sections for various gaugino production modes at the 14 TeV LHC. We have assumed the gaugino GUT relation, $M_3 : M_2 : M_1 = 7 : 2 : 1$, at the EW scale and plotted the cross sections as functions of M_2 (and $m_{\tilde{g}} \simeq M_3 = 7M_2/2$). The other relevant parameters were fixed as $m_{\tilde{q}} = \mu = 3$ TeV and $\tan\beta = 10$.

One can see that the $\tilde{\chi}_2^0\tilde{\chi}_1^\pm$ and $\tilde{\chi}_1^+\tilde{\chi}_1^-$ production modes have substantial cross sections. Because of the large mass hierarchy in the gaugino GUT relation, the $\tilde{g}\tilde{g}$ cross section drops much faster than the EW gaugino production cross sections as M_2 increases. Due to this effect, $\tilde{\chi}_2^0\tilde{\chi}_1^\pm$ and $\tilde{\chi}_1^+\tilde{\chi}_1^-$ production dominate over $\tilde{g}\tilde{g}$ production for $M_2 \gtrsim 350$ GeV.

The EW gaugino production modes other than $\tilde{\chi}_2^0\tilde{\chi}_1^\pm$ and $\tilde{\chi}_1^+\tilde{\chi}_1^-$ have cross sections which are a few orders of magnitude smaller. This is due to the fact that these production modes contain at least one bino state or two \tilde{W}^0 states in the large μ limit, and there exists no gaugino-gaugino-gauge boson couplings for those states.

As can be seen from figures 2 and 3, the $\tilde{\chi}_2^0\tilde{\chi}_1^\pm$ cross section is more than two times larger than the $\tilde{\chi}_1^+\tilde{\chi}_1^-$ cross section. This is mainly because $\tilde{\chi}_2^0\tilde{\chi}_1^\pm$ contains two distinctive modes: $\tilde{\chi}_2^0\tilde{\chi}_1^+$ with W^+ exchange and $\tilde{\chi}_2^0\tilde{\chi}_1^-$ with W^- exchange. It is therefore more beneficial to target the $\tilde{\chi}_2^0\tilde{\chi}_1^\pm$ production mode in the EW gaugino searches.

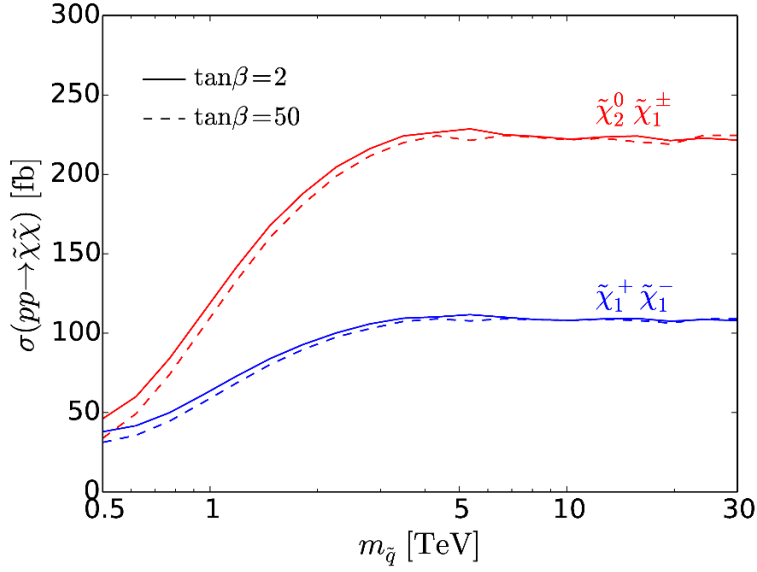


Figure 2. The NLO production cross sections for the $\tilde{\chi}_1^0 \tilde{\chi}_1^\pm$ and $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ modes at the 14 TeV LHC as functions of the squark mass. The cross sections have been calculated using **Prospino 2.1** [42, 43] with all the charges summed. We have set $\mu = m_{\tilde{q}}$ and $M_2 = 350$ GeV and $M_1 = 100$ GeV. The solid and dashed curves correspond to $\tan\beta = 2$ and 50, respectively.

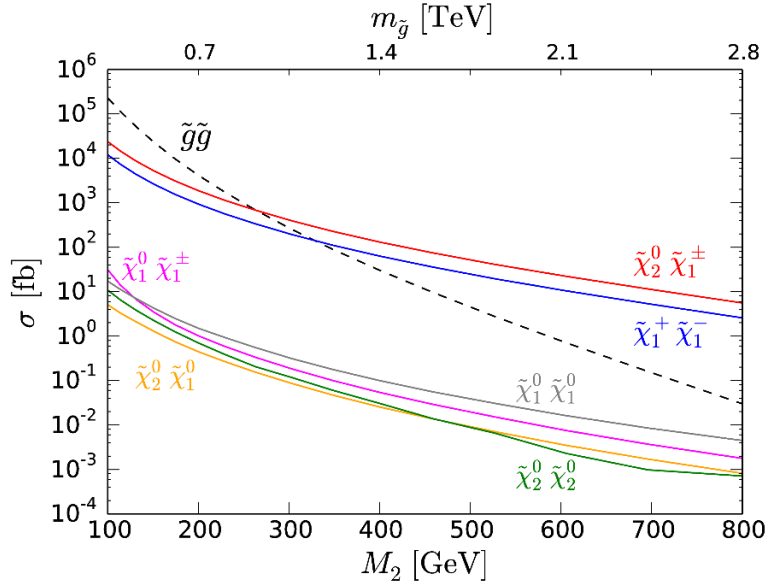


Figure 3. The NLO cross sections for various gaugino production modes at the 14 TeV LHC as functions of M_2 (and $m_{\tilde{g}} \simeq M_3 = 7M_2/2$). We have assumed the gaugino GUT relation, $M_3 : M_2 : M_1 = 7 : 2 : 1$, at the EW scale. The other relevant parameters were fixed as $m_{\tilde{q}} = \mu = 3$ TeV and $\tan\beta = 10$.

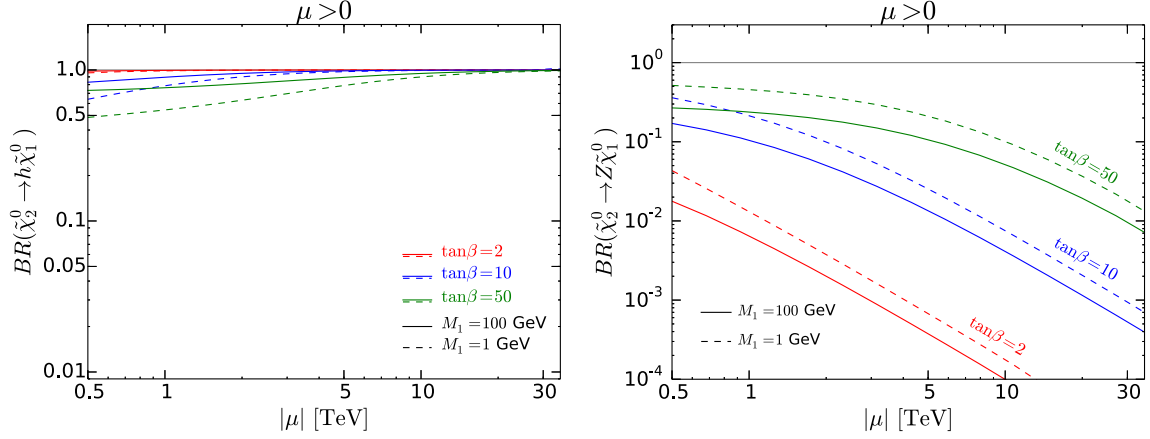


Figure 4. The branching ratios of $\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0$ (left) and $\tilde{\chi}_2^0 \rightarrow Z\tilde{\chi}_1^0$ (right) modes as functions of $|\mu|$ in the $\mu > 0$ case. M_2 is taken to be 350 GeV and $m_h = 125.5$ GeV. We have fixed $M_2 = 350$ GeV but we show variations of $\tan\beta$ and M_1 as $\tan\beta = 2$ (red), 10 (blue), 50 (green) and $M_1 = 100$ GeV (solid), 1 GeV (dashed).

2.3 The branching ratios

If scalar fermions and the MSSM Higgs bosons (other than the SM-like one) are heavier than the $\tilde{\chi}_1^\pm$ and $\tilde{\chi}_2^0$, these gaugino states decay predominantly into $\tilde{\chi}_1^0$ and SM bosons, W^\pm , Z and h , if the decays are kinematically allowed. In this case, $\tilde{\chi}_1^\pm$ exclusively decays into W^\pm and $\tilde{\chi}_1^0$ with $\text{BR} \sim 100\%$. On the other hand, $\tilde{\chi}_2^0$ has two possible decay modes: $\tilde{\chi}_2^0 \rightarrow Z\tilde{\chi}_1^0$ and $\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0$. The decay rates of these modes are determined by the $\tilde{\chi}_2^0\tilde{\chi}_1^0 Z/h$ couplings, up to the phase space factor and the polarisation effect. In the limit of large $|\mu|$ and heavy MSSM Higgs bosons, the $\tilde{\chi}_2^0\tilde{\chi}_1^0 Z/h$ couplings in the CP-conserving case are given by [44, 45]

$$\begin{aligned} |C_{\tilde{\chi}_1^0\tilde{\chi}_2^0 Z}| &\simeq \frac{e}{2} \frac{m_Z^2}{|\mu|^2}, \\ |C_{\tilde{\chi}_1^0\tilde{\chi}_2^0 h}| &\simeq \frac{e}{2} \frac{m_Z}{|\mu|} \left| 2\sin 2\beta + \frac{M_1 + M_2}{\mu} \right|, \end{aligned} \quad (2.2)$$

where e is the electric charge ($\alpha_{em} = e^2/(4\pi)$).

Figure 4 shows the branching ratios of $\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0$ (left) and $\tilde{\chi}_2^0 \rightarrow Z\tilde{\chi}_1^0$ (right) modes as functions of $|\mu|$ in the $\mu > 0$ case. M_2 has been fixed to $M_2 = 350$ GeV but $\tan\beta$ and M_1 are varied as $\tan\beta = 2$ (red), 10 (blue), 50 (green) and $M_1 = 100$ GeV (solid), 1 GeV (dashed). Here and throughout the paper, we have explicitly set $m_h = 125.5$ GeV. This condition can be always realised by tuning the stop mass, which has no effect on the EW gaugino sector, and hence on our phenomenological analysis. The branching ratios were calculated using SUSY-HIT [46].

One can see that the $\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0$ mode is enhanced, whilst the $\tilde{\chi}_2^0 \rightarrow Z\tilde{\chi}_1^0$ mode is suppressed as $|\mu|$ increases. This is due to the $\tilde{\chi}_1^0\tilde{\chi}_2^0 Z$ coupling having the extra m_Z/μ suppression factor compared to the $\tilde{\chi}_1^0\tilde{\chi}_2^0 h$ coupling, as seen in eq. (2.2). In the $|\mu| \gtrsim 500$ GeV region, the $\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0$ mode has $\text{BR} \gtrsim 60\%$ and dominates the $\tilde{\chi}_2^0$ decay, apart from

the $\tan\beta = 50$, $M_1 = 1$ GeV case. For moderate values of μ , $0.5 \lesssim \mu/\text{TeV} \lesssim 3$, the factor $|2\sin 2\beta + (M_1 + M_2)/\mu|$ in $C_{\tilde{\chi}_1^0 \tilde{\chi}_2^0 h}$ is important in the competition between the $\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0$ and $\tilde{\chi}_2^0 \rightarrow Z\tilde{\chi}_1^0$ modes and at $\mu \sim 1$ TeV, $\tan\beta \sim 50$, $M_1 \sim 1$ GeV, $BR(\tilde{\chi}_2^0 \rightarrow Z\tilde{\chi}_1^0)$ can be as large as $BR(\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0)$. However, in the large $|\mu|$ limit $BR(\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0)$ approaches 100 % independently of $\tan\beta$ and M_1 as long as the phase space is open.

Figure 5 is equivalent to figure 4, with μ instead set to $\mu < 0$. One can see that $BR(\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0)$ becomes zero at a particular $|\mu|$ value depending on $\tan\beta$ and M_1 . This is due to the cancellation between the two terms in the $|2\sin 2\beta + (M_1 + M_2)/\mu|$ factor in the $\tilde{\chi}_1^0 \tilde{\chi}_2^0 h$ coupling. As can be seen, this cancellation occurs at $\mu \sim -1$ TeV for $\tan\beta \sim 10$ and $\mu \sim -5$ TeV for $\tan\beta \sim 50$. As $|\mu|$ increases, $BR(\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0)$ quickly approaches 100 % following the cancellation. For $|\mu| \gtrsim 10$ TeV, the $\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0$ mode dominates over the $\tilde{\chi}_2^0$ decay, independently of $\tan\beta$ and M_1 .

To summarise, we have demonstrated that in the scenarios with large $m_{\tilde{q}}$ and $|\mu|$, $\tilde{\chi}_2^0$ and $\tilde{\chi}_1^\pm$ become wino-like gauginos with $m_{\tilde{\chi}_2^0} \simeq m_{\tilde{\chi}_1^\pm} \simeq M_2$ and $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm$ has the largest cross section among the EW gaugino production modes. We also argued that in such scenarios $\tilde{\chi}_1^\pm$ predominantly decays into W^\pm and $\tilde{\chi}_1^0$ and the $\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0$ mode typically dominates the $\tilde{\chi}_2^0$ decay. These arguments provide a strong motivation to study the $pp \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow (h\tilde{\chi}_1^0)(W^\pm \tilde{\chi}_1^0)$ mode in the EW gaugino searches in the scenarios with large $m_{\tilde{q}}$ and $|\mu|$.

In the following sections, we study the $pp \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow (h\tilde{\chi}_1^0)(W^\pm \tilde{\chi}_1^0)$ channel using the $W \rightarrow \tau/\ell, \nu$ plus $h \rightarrow \tau\tau$ and $h \rightarrow WW \rightarrow (\tau/\ell, \nu)(\tau/\ell, \nu)$ channel. We set $m_{\tilde{q}} = \mu = m_A = 3$ TeV, $\tan\beta = 10$ throughout. This leads to $BR(\tilde{\chi}_1^\pm \rightarrow W^\pm \tilde{\chi}_1^0) \simeq BR(\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0) \simeq 100$ %. With this parameter choice, the lightest CP-even Higgs becomes SM-like and we use the same branching ratios as those for the SM Higgs boson. Although the above parameter set is motivated by the heavy scalar scenario, our analysis can easily be recast onto other SUSY scenarios. Changing the above parameters may modify the $pp \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^\pm$ cross section and the $\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0$ branching ratio significantly but does not alter the signal efficiencies for the signal regions defined in the next section. The discovery reach and exclusion limit for a different set of parameters can therefore be obtained by rescaling the cross section and branching ratio accordingly. Moreover, the calculated signal efficiencies can also be used for a larger class of models as far as the $\tilde{N}\tilde{C}^\pm \rightarrow (h\chi)(W^\pm \chi)$ topology is concerned, as mentioned in subsection 2.1. Neglecting a finite width effect and spin correlations, the signal efficiencies will be very similar between $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow (h\tilde{\chi}_1^0)(W^\pm \tilde{\chi}_1^0)$ and $\tilde{N}\tilde{C}^\pm \rightarrow (h\chi)(W^\pm \chi)$ at $(m_{\tilde{\chi}_2^0} = m_{\tilde{\chi}_1^\pm}, m_{\tilde{\chi}_1^0}) = (m_{\tilde{N}} = m_{\tilde{C}}, m_\chi)$. We explain this point in more detail in appendix F and provide the necessary information to perform such a re-analysis.

3 Simulation and analysis

3.1 Monte Carlo simulation

The SUSY $pp \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow (h\tilde{\chi}_1^0)(W^\pm \tilde{\chi}_1^0)$ signals were generated using the HERWIG++ general-purpose event generator [47–49] via SUSY Les Houches Accord files used as input for the parameter points, according to the assumptions outlined in the previous section. The signal cross sections were scaled to the next-to-leading order cross sections using results

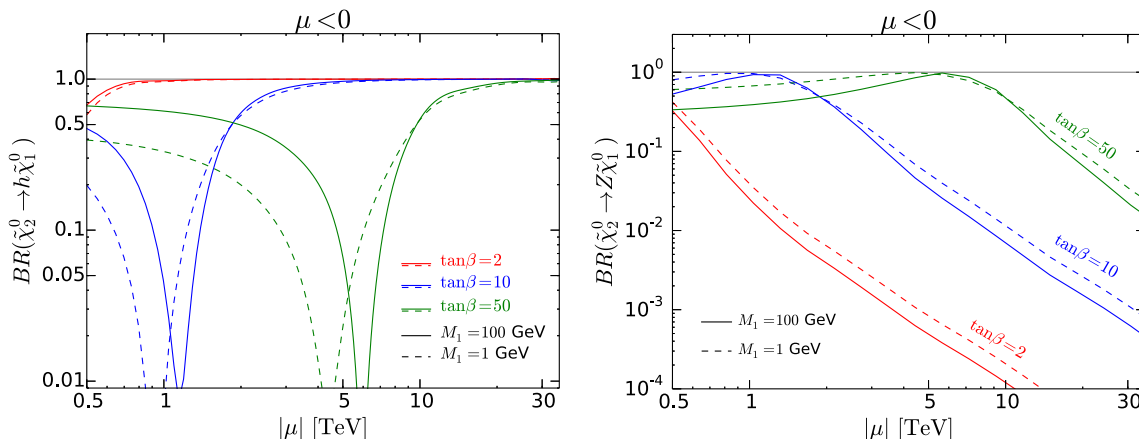


Figure 5. Equivalent plots to figure 4 but with $\mu < 0$.

obtained from Prospino 2.1. The hV , $t\bar{t}$, $t\bar{t}h$ and WZ backgrounds were also generated internally in HERWIG++ at leading order. The Z +jets and W +jets backgrounds were generated using the parton-level matrix element generator AlpGen and merged with the HERWIG++ parton shower using the MLM method [50–52]. The generator-level cuts on the V +jets backgrounds were taken to be $p_{Tj,\min} = 15$ GeV, $\eta_{j,\max} = 3.0$, $\Delta R_{j,\min} = 0.2$ with $m_{\ell\ell} \in (15, 160)$ GeV (or $m_{\tau,\tau}$) for $V = Z$. For the Z +jets case we considered matrix elements with one extra parton merged to the shower, whereas for the W +jets case we considered matrix elements with two partons merged to the shower.

For the signal we allowed the W to decay to all lepton flavours, including taus. Likewise, for the backgrounds we consider all of the leptonic decays of the W and Z , to muons, electrons or taus. We consider the Monte Carlo samples of the Z and W backgrounds going to electrons or muons separately from those going to taus, as they would have different amounts of missing energy, leptons and jets. The Higgs boson was allowed to decay to $\tau^+\tau^-$ or W^+W^- with subsequent decay of the W bosons to $e\nu_e$, $\mu\nu_\mu$ and $\tau\nu_\tau$.

In all cases of signal and background the full parton shower, hadronization and the underlying event [53] were included.³ All the runs have been generated using the MSTW2008nlo 68% PDF set. We note that we do not consider pure QCD-initiated backgrounds since these are expected to be negligible in the high-missing transverse momentum regime, particularly in conjunction with the existence of isolated leptons or τ s.

We define a SUSY benchmark point **C350-100**, with parameters

$$M_2 = m_{\tilde{\chi}_1^\pm} = m_{\tilde{\chi}_2^0} = 350 \text{ GeV}, \quad M_1 = m_{\tilde{\chi}_1^0} = 100 \text{ GeV}. \quad (3.1)$$

This point will be used as an example to demonstrate the effect of cuts and provide a typical point to aid the development of the strategy for discriminating the signal against the various backgrounds.

³The event generator dependence is beyond the scope of this phenomenological study. Moreover, we note that do not include a description of pile-up events. These should be considered in detail in a full experimental simulation.

3.2 Tau identification

3.2.1 Tau lepton decay modes

The study of final states containing hadronically decaying τ leptons is an important and growing part of the LHC's physics program. The τ lepton has a multitude of decay modes, which we may split these into two categories: ‘leptonic’, if the visible decay products contain a single lepton, and ‘hadronic’, if there are one or three charged hadrons present. We label the corresponding modes τ_ℓ and τ_h respectively. The hadronic modes are also categorised as ‘1-prong’ and ‘3-prong’, according to the number of charged particles involved in the decay.⁴ The label ‘ ℓ ’ here and elsewhere implies an electron or a muon. The branching ratios for these modes are:⁵

- leptonic: $\text{BR}(\tau \rightarrow \tau_\ell) \sim 0.35$.
- hadronic: $\text{BR}(\tau \rightarrow \tau_h) \sim 0.625$.

These imply, for a Higgs boson decay to $\tau^+\tau^-$:

- $\text{BR}(h \rightarrow \tau_h\tau_h) \sim 0.39 \times \text{BR}(h \rightarrow \tau^+\tau^-)$.
- $\text{BR}(h \rightarrow \tau_h\tau_\ell) \sim 0.44 \times \text{BR}(h \rightarrow \tau^+\tau^-)$.
- $\text{BR}(h \rightarrow \tau_\ell\tau_\ell) \sim 0.12 \times \text{BR}(h \rightarrow \tau^+\tau^-)$.

3.2.2 Hadronic tau identification algorithm

Both ATLAS [55] and CMS [56] employ reconstruction and identification algorithms, used to identify hadronically decaying τ leptons and reject various backgrounds. Here, we do not attempt to reproduce either of the ATLAS or CMS algorithms exactly, but instead use elements from both resulting in an algorithm that we expect performs in an equivalent way. We also borrow elements from [57], which examines di- τ tagging in the boosted regime.⁶ The resulting algorithm is expected to provide conservative hadronic τ -tagging results, and could be improved substantially via the use of boosted decision trees (BDT) or other advanced multivariate methods. Since we will not employ simulation of detector effects in the present analysis, we focus on a simple cut-based algorithm for simplicity.

The first part of the basic algorithm for hadronic τ identification proceeds as follows:

- Reconstruct jets with $R = 0.5$ using the Cambridge/Aachen jet algorithm as implemented in **FastJet** [59]. An individual jet is then investigated for constituent hadronic tracks.⁷

⁴If one of the taus undergoes a 3-prong decay, one may improve the analysis significantly using the information of the secondary vertex of the 3-prong tau decay [54]. This requires a dedicated study and we do not use the secondary vertex information in this paper.

⁵These do not add up to 100%, since we are only considering the dominant 1-prong and 3-prong decay modes.

⁶For di- τ tagging in Higgs searches, see also ref. [58].

⁷Usage of the word “track” here and elsewhere in this article implies “charged particle”.

- Consider a track to be a ‘seed’ if it is the hardest track in the jet, has $p_T > 5 \text{ GeV}$ and is within $\Delta R = 0.1$ of the jet axis.
- If such a track is found, one defines inner and outer cones around it. We use $R_{\text{in}} = 0.2$ and $R_{\text{out}} = 0.4$ respectively.
- Require *no* photons with $p_T > 2 \text{ GeV}$ and *no* charged tracks with $p_T > 1 \text{ GeV}$ to lie within the defined annulus between R_{in} and R_{out} .

The basic part of the algorithm itself does not provide satisfactory rejection against the QCD jet background to hadronically decaying τ leptons. If a jet satisfies all the above criteria, then the following variables are constructed:

- ΔR_{max} : the distance to the track furthest away from the jet axis.
- f_{core} : the fraction of the total jet energy contained in the centre-most cone defined by $\Delta R < 0.1$.

These variables provide strong discriminating power against QCD jets [55, 60].⁸ To perform the rejection of QCD jets, here we apply the following cuts:

- $\Delta R_{\text{max}} < 0.05$.
- $f_{\text{core}} > 0.95$.

In figure 6 we show the variables ΔR_{max} and f_{core} , constructed for hadronic jets for a signal benchmark point **C350-100** and the W +jets background.⁹ Only jets with $p_T > 20 \text{ GeV}$ and $|\eta| < 2.5$ were considered. In figure 7 we show the efficiency of τ identification versus the transverse momentum of the jet in question, $p_{T,\text{jet}}$, obtained by the procedure outlined in this section. For the signal, the efficiency was defined for the identification of ‘true’ τ jets, defined to be those closest to the visible τ decay products taken from the Monte Carlo truth. For the W +jets background, the efficiency was defined with respect to any jet. The efficiency for the SUSY benchmark point **C350-100** varies from around 50% in the $p_{T,\text{jet}}$ region of $20 - 300 \text{ GeV}$ and then drops down to $\sim 20\%$ at around $p_{T,\text{jet}} \sim 400 \text{ GeV}$. For the W +jets background the efficiency starts off at $\sim 1\%$ at $p_{T,\text{jet}} \sim 20 \text{ GeV}$ and then rises to an efficiency of $2 - 3\%$, more or less constant up to $p_{T,\text{jet}} \sim 500 \text{ GeV}$.

3.3 Analysis

In order not to limit the scope of our phenomenological analysis, we will assume that a multi-object trigger can conceivably be designed, involving multiple leptons or hadronic tau jets with $p_T \gtrsim 20 \text{ GeV}$, in conjunction with missing transverse energy $\cancel{E}_T \gtrsim 100 \text{ GeV}$ [61–63]. Such a trigger actually appears in the ATLAS study of ref. [64], where in section 3.4, pre-selection criteria of $p_T(\text{lep}) > 10 \text{ GeV}$, $p_T(\text{tau}) > 20 \text{ GeV}$ and $\cancel{E}_T > 100 \text{ GeV}$

⁸More variables have been employed by the experimental collaborations, but we found that the two that we consider are sufficient at this level of simulation.

⁹Two-dimensional distributions of the two variables for the two samples that we have investigated can be seen in appendix A.

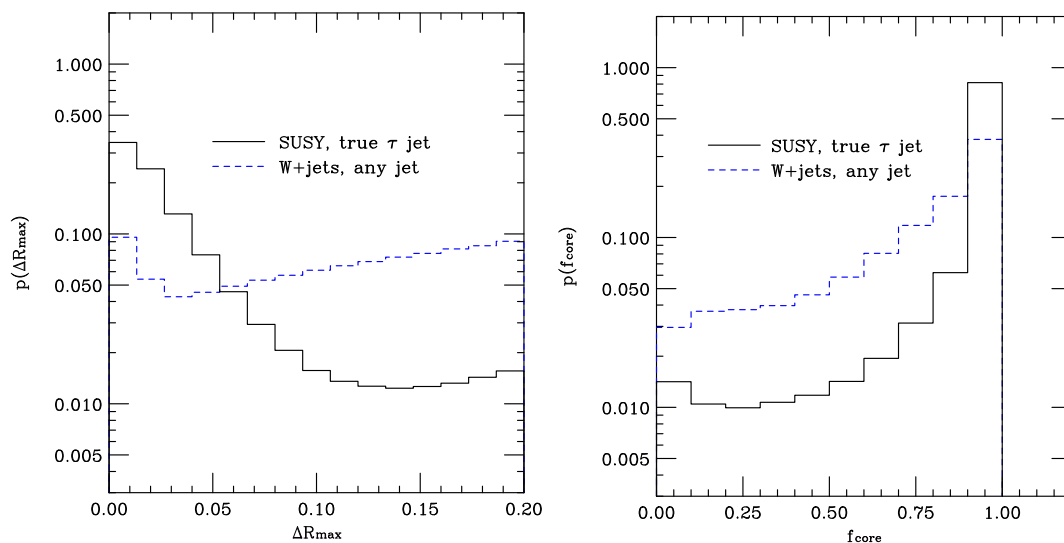


Figure 6. Distributions of the variables used for discrimination of the jets originating from τ leptons and those from QCD, for the SUSY benchmark point **C350-100** and the W +jets background ($W \rightarrow e\nu_e/\mu\nu_\mu$).

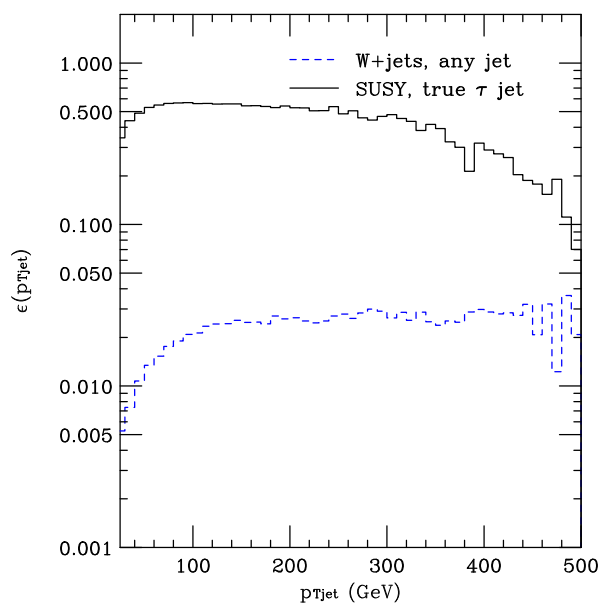


Figure 7. The efficiency of tagging a jet as a τ -jet for the SUSY benchmark point **C350-100** and the W +jets background (with $W \rightarrow e\nu_e/\mu\nu_\mu$). For **C350-100**, the efficiency was defined for the identification of ‘true’ τ jets, defined to be those closest to the visible τ decay products taken from the Monte Carlo truth. For the W +jets, the efficiency was defined with respect to any jet. Jets of $p_T > 20$ GeV and $|\eta| < 2.5$ are considered in both cases.

are imposed for a similar Wh -mediated SUSY search at the high-luminosity LHC. Our selection criteria are essentially tighter than their pre-selection criteria. We investigate the effect of increasing the p_T thresholds for leptons and taus below.

We define the first level of the analysis for discriminating the signal against the various backgrounds as follows:

1. Particles of $p_T > 0.5 \text{ GeV}$ and $|\eta| < 5.0$ are considered.
2. If isolated leptons with $p_T > 20 \text{ GeV}$ are found, they are placed in a separate list, and removed from the list of particles. An isolated lepton is defined as either: having $\sum_i p_{T,i}$ less than 20% of its transverse momentum around a cone of $\Delta R = 0.4$ around it, or as a lepton that contains no photons with $p_T > 2 \text{ GeV}$ and no tracks with $p_T > 1 \text{ GeV}$ in the annulus $\Delta R = (0.2, 0.4)$ around it.¹⁰
3. Jet finding is performed on the list of remaining particles, using **FastJet** and the Cambridge/Aachen jet algorithm, with parameter $R = 0.5$. Jets of $p_T > 20 \text{ GeV}$ are accepted.
4. Tagging of τ -jets is performed as described in section 3.2.
5. Only events with a total number of isolated leptons, $n_{\ell,\text{iso}}$, and τ -tagged jets, $n_{\tau,\text{tag}}$, equal to 3 are accepted: i.e. we require $n_{\tau,\text{tag}} + n_{\ell,\text{iso}} = 3$. A hypothesis is then performed to match the topology of the SUSY events. The hypotheses vary according to the number of isolated leptons and τ -tagged jets and are listed in detail in Table 2.
6. Several variables are calculated and are passed through to the second level of analysis.

Steps 1-5 are what we define as the ‘basic’ analysis, and these will provide what we call σ_{basic} (table 4). The variables calculated in step 6 and used for further discrimination in the second-level analysis are: the transverse momentum of the di- τ -tagged system, $p_{T,\tau\tau}$, the distance between the τ -tagged jets, $\Delta R_{\tau,\tau}$, the distance between the di- τ -tagged system and the lepton, $\Delta R_{\tau\tau,\ell}$, the missing transverse energy, \cancel{p}_T and the variable M_{min} , which is sharply peaked at low values for the WZ background and broadly falls off for the signal, defined in appendix B. The variables are outlined in table 3. There we provide an *example* set of cuts, applied to the SUSY benchmark point **C350-100**, found to give a significance of $\sim 2.5\sigma$ for an integrated luminosity of 100 fb^{-1} at 14 TeV. To investigate the effect of higher p_T thresholds, we set for each isolated lepton, $p_T(\text{lept}) > 30 \text{ GeV}$, and for each hadronic tau, $p_T(\text{tau}) > 40 \text{ GeV}$. This reduces the significance to $\sim 2\sigma$ and verifies that our analysis is not substantially affected by higher thresholds that aim to take into account current triggering strategies.

For completeness, we show in table 4 the resulting cross sections after applying the analysis on the SUSY benchmark point and the different backgrounds for this example. In the final stage of the analysis the $n_{\tau,\text{tag}} = 3$ channel was excluded, since it was found to reduce significance by allowing more background. Note that this set of cuts will constitute ‘signal region 1’ of our full analysis.

¹⁰We apply two different criteria to take into account the possibility of radiation from the core lepton.

$n_{\tau,\text{tag}}$	$n_{\ell,\text{iso}}$	real signal channels	hypothesis
3	0	$(h \rightarrow \tau_h \tau_h, W \rightarrow \tau \nu)$	assign hardest two to h .
2	1	$(h \rightarrow \tau_h \tau_\ell, W \rightarrow \tau_h \nu_\tau),$ $(h \rightarrow \tau_h \tau_h, W \rightarrow \ell \nu_\ell)$	assign hardest two to h .
1	2	$(h \rightarrow \tau_h \tau_\ell, W \rightarrow \ell \nu_\ell),$ $(h \rightarrow \tau_\ell \tau_\ell, W \rightarrow \tau_h \nu_\tau)$	if leptons are same sign, assign highest- p_T to h along with the τ -tagged jet. Otherwise: assign any two highest- p_T to h .
0	3	$(h \rightarrow \tau_\ell \tau_\ell, W \rightarrow \ell \nu_\ell)$	If all leptons are the same sign, reject the event. Otherwise: pair two highest- p_T of opposite sign as the h .

Table 2. The hypotheses applied for the reconstruction of the Supersymmetric topology as described in the main text. The different hypotheses are given according to the number of τ -tagged jets, $n_{\tau,\text{tag}}$, and the number of isolated leptons $n_{\ell,\text{iso}}$. In the final stage of the analysis, the $n_{\tau,\text{tag}} = 3$ was found to reduce significance and was not considered.

Details of how the initial cross sections for the signal and background are calculated are given in appendix C. We note that in the case of the Z +jets and W +jets samples, we obtained $N_{\text{cuts}} = 0$ events after all cuts.¹¹ To provide an estimate of the cross section, we assume that the Poisson distribution has mean number of events $\lambda = 3$ and use this as an upper bound to estimate the resulting cross sections. The probability of having a Poisson-distributed sample with mean $\lambda > 3$, given that zero events have been observed, is $\simeq 0.05$. It is useful to mention at this point that we do not apply a K -factor to the Z +jets or W +jets cross sections. The induced uncertainty due to this omission can be absorbed in the systematic uncertainty due to lack or low number of events in the final Monte Carlo samples. Nevertheless, since conservative estimates for these backgrounds have been assumed, K -factors of ~ 2 would not have a significant impact to the main conclusions of our analysis.

3.4 Signal regions

To perform a scan of the supersymmetric parameter space, we define signal regions, with cuts that aim to bring out the different qualities of the defined variables. These signal regions are shown in table 5, for the variables defined in table 3. All the signal regions exclude the $n_{\tau,\text{tag}} = 3$ channel, since it was found to reduce significance.

4 Results

We performed the analysis on the M_2 - M_1 plane, according to the cuts defined in the signal regions in table 5 at integrated luminosities of 100 fb^{-1} , 300 fb^{-1} and 3000 fb^{-1} . We show

¹¹For the $Z(\rightarrow \tau^+ \tau^-)$ +jets, this depends on how large the missing transverse momentum cut imposed is.

variable	definition	benchmark point cut (\equiv signal region 1)
\not{p}_T	missing transverse momentum	$> 95 \text{ GeV}$
M_{\min}	appendix B	$> 235 \text{ GeV}$
$p_{T,\tau\tau}$	di- τ -tagged jet p_T	$> 20 \text{ GeV}$
$\Delta R_{\tau,\tau}$	distance between τ -tagged jets	$\in (0.1, 2.9)$
$\Delta R_{\tau\tau,\ell}$	distance between di- τ -tagged jet system and lepton	$\in (0.1, 2.6)$

Table 3. The variables used for further discrimination after the basic part of the analysis is applied to the signal and backgrounds.

sample	σ_{initial} (fb)	σ_{basic} (fb)	σ_{cuts} (fb)
SUSY C350-100	5.7	0.658	0.152
WZ	767	85.734	0.079
$W(\rightarrow \ell\nu_\ell)+\text{jets}$	$\sim 600 \times 10^3$	61.974	$\lesssim 0.055$
$W(\rightarrow \tau\nu_\tau)+\text{jets}$	$\sim 300 \times 10^3$	7.591	$\lesssim 0.052$
hV	443	5.071	0.037
$t\bar{t}h$	3.4	0.147	0.008
$t\bar{t}$	8600	14.876	0.005
$Z(\rightarrow \ell\ell)+\text{jets}$	$\sim 600 \times 10^3$	1659	$\lesssim 0.029$
$Z(\rightarrow \tau\tau)+\text{jets}$	$\sim 300 \times 10^3$	52.762	0.047

Table 4. The effect of the cuts on the SUSY benchmark point **C350-100** and the relevant backgrounds. The ‘initial’ cross section calculations are presented in appendix C. The ‘basic’ cross sections are obtained after applying the procedure outlined in steps 1-5 in section 3.3.

variable	SR1	SR2	SR3	SR4	SR5	SR6	SR7
\not{p}_T	95 GeV	120 GeV	100 GeV	90 GeV	90 GeV	150 GeV	90 GeV
M_{\min}	235 GeV	270 GeV	220 GeV	220 GeV	300 GeV	240 GeV	200 GeV
$p_{T,\tau\tau}$	20 GeV	80 GeV	20 GeV	50 GeV	20 GeV	20 GeV	20 GeV
$\Delta R_{\tau,\tau}$	(0.1, 2.9)	(0.1, 2.9)	(0.1, 2.9)	(0.1, 2.9)	(0.1, 2.9)	(0.1, 2.9)	(0.1, 2.9)
$\Delta R_{\tau\tau,\ell}$	(0.1, 2.6)	(0.1, 2.5)	(0.1, 2.6)	(0.1, 2.6)	(0.1, 2.6)	(0.1, 2.6)	(0.1, 2.6)

Table 5. The cuts for the different signal regions (SR) used in the analysis.

the resulting envelope of significances in figure 8, where the solid curves show the 3σ evidence region, whereas the dashed curves show the 5σ discovery region. We also show in figure 9, the expected exclusion region at 2σ (solid) and 3σ (dashed). For completeness, we show the corresponding overlapping signal regions in appendix F. There, we also provide the total cross sections for the backgrounds after cuts given by the different signal regions. These can be used to infer constraints in explicit SUSY models that contain the specific decay chain we are considering.

The analysis can yield a low number of events for both signal and background, of $\mathcal{O}(10)$, and for the calculation of significance we used the Poisson distribution to calculate the p -values. These were subsequently converted to the corresponding Gaussian standard deviations. Details of the procedure are provided in appendix D, with supplementary material in appendix E.

Although the authors of ref. [34] have not performed an equivalent parameter-space scan over M_1 - M_2 , and the details of the chosen parameters differ from the ones presented in this article, it is still interesting to compare with the potential of the final state in which the channel $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow (h \tilde{\chi}_1^0)(W^\pm \tilde{\chi}_1^0)$ involves leptonic W decays and Higgs boson decays to $b\bar{b}$. There, the authors have found that it is possible to discover a signal of the process at the $\sim 5\sigma$ level at $\sim 100 \text{ fb}^{-1}$ of luminosity, for points for which $M_2 \sim 265 - 390 \text{ GeV}$ and $M_1 \sim 133 - 198 \text{ GeV}$. Indeed, our analysis is competitive with this result, with such points falling somewhere between the 3σ and 5σ discovery regions at 100 fb^{-1} and 300 fb^{-1} , as demonstrated by the black and red curves respectively, in figure 8. This indicates that this channel is as important as the final state with $h \rightarrow b\bar{b}$, or at least complementary.

5 Conclusions

We have presented a phenomenological analysis of the channel $\tilde{\chi}_2^0 \tilde{\chi}_1^\pm \rightarrow (h \tilde{\chi}_1^0)(W^\pm \tilde{\chi}_1^0)$ using the $W \rightarrow \ell \nu_\ell / \tau \nu_\tau$ and Higgs boson channels ($h \rightarrow \tau^+ \tau^-$ and $h \rightarrow W^+ W^- \rightarrow \text{leptons}$) at the LHC. Such channels are common in many concrete SUSY models where the predictions include $\tilde{\chi}_2^0$ and $\tilde{\chi}_1^\pm$ that predominantly decay into h and W , respectively.

Our analysis has included detailed hadron-level simulation of the relevant dominant backgrounds, including the effects of the underlying event. Hadronic τ identification was modelled at hadron level with a custom-made algorithm based on the ones employed by both the ATLAS and CMS experiments. We have employed a cut-based analysis on several variables that bring out the properties of the signal against those of the backgrounds. Specifically, we have constructed a mass variable, M_{\min} , which is sharply peaked at low value for the WZ background and broadly falls off for the signal.

Consequently we have demonstrated the potential for discovering or constraining the SUSY parameter space in the M_2 - M_1 plane at integrated luminosities of 100 fb^{-1} , 300 fb^{-1} and 3000 fb^{-1} , collected at a 14 TeV proton-proton centre-of-mass energy. The 5σ discovery potential of our analysis reaches up to $M_2 \simeq 350 \text{ GeV}$ with $M_1 \lesssim 100 \text{ GeV}$ at the 14 TeV LHC with 300 fb^{-1} . This implies that a future e^+e^- collider with $\sqrt{s} = 1 \text{ TeV}$ can play indispensable role to cover $M_2 < 500 \text{ GeV}$ region. A large part of this region can also

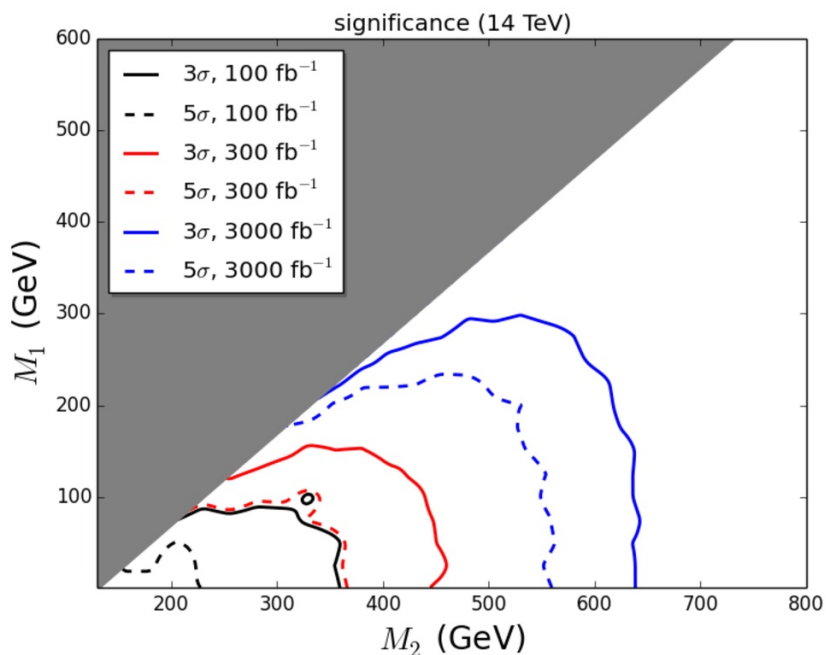


Figure 8. The significance envelope on the M_2 - M_1 plane obtained for the signal regions defined in table 5 at integrated luminosities of 100 fb^{-1} (black) or 300 fb^{-1} (red). The solid curves show the 3σ evidence region, whereas the dashed curves show the 5σ discovery region.

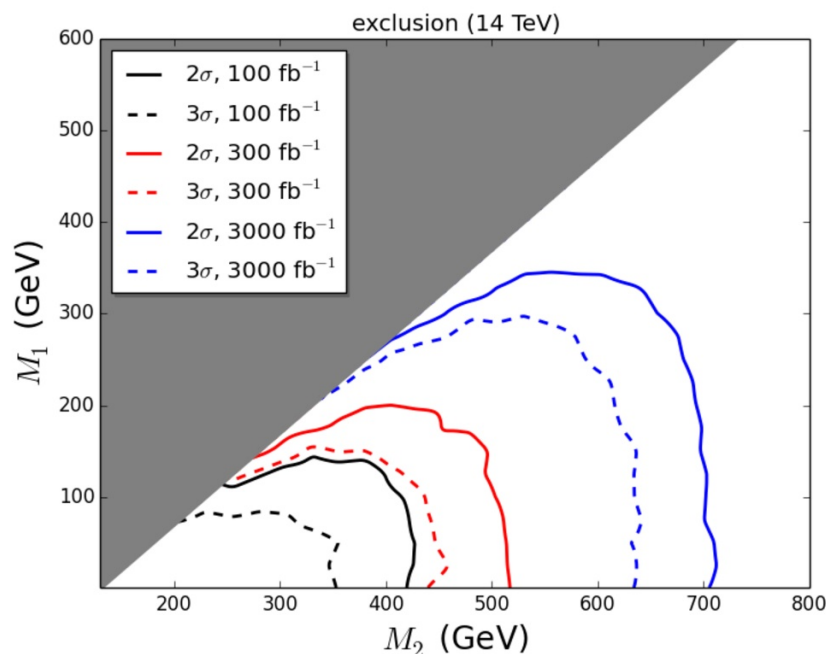


Figure 9. The exclusion envelope on the M_2 - M_1 plane obtained for the signal regions defined in table 5 at integrated luminosities of 100 fb^{-1} (black) or 300 fb^{-1} (red). The solid curves show the 2σ exclusion boundary, whereas the dashed curves show the 3σ boundary.

be covered by the 14 TeV High Luminosity LHC with 3000 fb^{-1} , which has a discovery potential in the $M_2 \lesssim 550 \text{ GeV}$, $M_1 \lesssim 200 \text{ GeV}$ region.

This work serves a first study of making use of $h \rightarrow \tau\tau$ mode in the chargino-neutralino searches in future runs of the LHC. We thus recommend further examination of this channel by experimental collaborations, including the effects of full detector simulation, τ -jet tagging and multi-variate analyses.

A Correlation between τ -tagging variables

We show in figure 10 the two-dimensional distributions of the two variables, ΔR_{max} and f_{core} , that we have employed for the discrimination of hadronic tau jets from QCD jets. The variables are correlated in both signal and background. However, it is evident that both variables are necessary for discrimination between the hadronic taus and the QCD jets. The figure has been normalised to integrate to unity over the range $(\Delta R_{\text{max}}, f_{\text{core}}) = (0.2, 1.0)$ and hence one can notice that the SUSY signal is more sharply peaked in the region $(\Delta R_{\text{max}} < 0.5, f_{\text{core}} > 0.95)$.

B Definition of the M_{min} variable

We define the M_{min} variable that we will use as a handle for rejecting non-SUSY backgrounds. Although the variable is designed to reject the WZ background, it can also potentially perform well against other backgrounds. There are three neutrinos in the final state: one coming from W decay, the other two from the τ lepton decays. The direction of the τ -neutrino is approximately collimated with respect to the original τ lepton direction due to the mass hierarchy, $m_Z \gg m_\tau$. With this approximation, the momenta of the τ lepton and the τ -neutrino can be parametrised as

$$\begin{aligned} p_{\tau^+} &= p_{\rho_1}/a, & p_{\tau^-} &= p_{\rho_2}/b, \\ p_{\nu_1} &= (1/a - 1)p_{\rho_1}, & p_{\nu_2} &= (1/b - 1)p_{\rho_2}, \end{aligned} \quad (\text{B.1})$$

where $p_{\rho_{1/2}}$ is the momentum of the visible decay products and: $0 < a(b) < 1$. Note that events that in the phenomenological analysis of this article that do not satisfy this condition on a and b are deemed ‘unphysical’ and rejected. Assuming the event topology in figure 11, the unknown neutrino momenta can be constrained by the mass shell conditions of the W and Z bosons and the missing momentum conditions.¹²

$$\begin{aligned} a, b, \mathbf{p}_\nu &: 5 \text{ unknowns} \\ m_Z, m_W, p_{\text{miss}}^x, p_{\text{miss}}^y &: 4 \text{ constraints} \end{aligned}$$

Since $(\# \text{ of unknown} - \# \text{ of constraints}) = 1$, we can parameterise the all neutrino momenta by a single parameter, θ .

¹²Vectors in bold typeset represent 3-vectors.

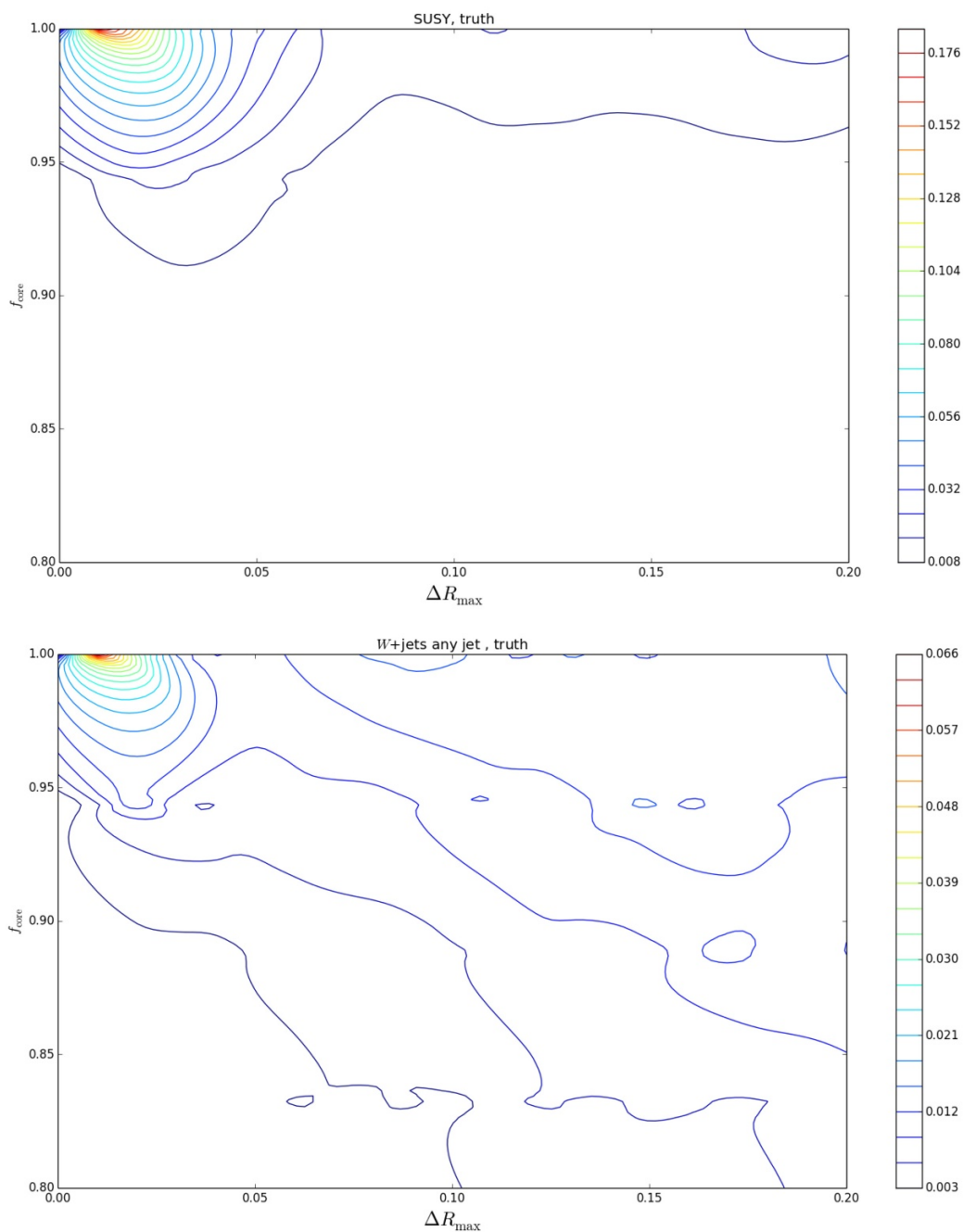


Figure 10. Two dimensional-distributions of the variables used for discrimination of the jets originating from τ leptons and those from QCD, ΔR_{\max} and f_{core} . top: SUSY benchmark point C350-100, bottom: W +jets background ($W \rightarrow e\nu_e/\mu\nu_\mu$).

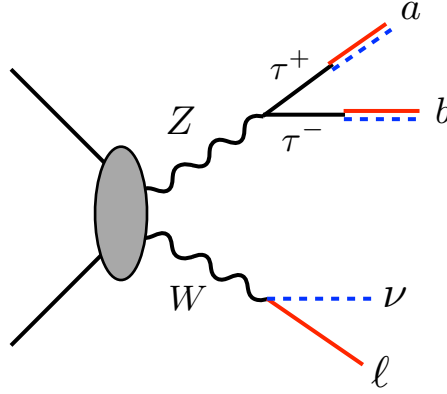


Figure 11. The WZ background topology considered in constructing the M_{\min} variable.

The mass-shell constraint for the Z boson gives

$$ab = \frac{2(p_{\rho_1} \cdot p_{\rho_2})}{m_Z^2}. \quad (\text{B.2})$$

By introducing $\theta \equiv \arctan\left(\frac{a}{b}\right)$, a and b can be written as

$$a = \sqrt{\frac{2(p_{\rho_1} \cdot p_{\rho_2})}{m_Z^2} \tan \theta}, \quad b = \sqrt{\frac{2(p_{\rho_1} \cdot p_{\rho_2})}{m_Z^2} \tan^{-1} \theta} \quad (\text{B.3})$$

The transverse components of the neutrino momentum are determined by

$$\mathbf{p}_\nu^T = \mathbf{p}_{\text{miss}}^T - (1/a - 1)\mathbf{p}_{\rho_1} - (1/b - 1)\mathbf{p}_{\rho_2}. \quad (\text{B.4})$$

The mass shell condition of W constrains the last unknown parameter p_ν^z as

$$p_\nu^{z\pm} = \frac{cp_\ell^z \pm \sqrt{E_\ell^2(c^2 - t_\ell^2 t_\nu^2)}}{t_\ell^2}, \quad (\text{B.5})$$

where $\mathbf{t}_{\ell/\nu} = \mathbf{p}_{\ell/\nu}^T$, $c = \mathbf{t}_\ell \cdot \mathbf{t}_\nu + m_W^2/2$. If eq. (B.5) yields complex solution, we simply take the real part [65, 66].

All the neutrino momenta are now parametrised by θ . We define the invariant mass of the system

$$M_{\text{inv}}^\pm(\theta) = \sqrt{[p_\ell + p_\nu^\pm(\theta) + p_{\tau^+}(\theta) + p_{\tau^-}(\theta)]^2}, \quad (\text{B.6})$$

where \pm corresponds to the discrete ambiguity in eq. (B.5). The variable M_{\min} is defined by the global minimum of the M_{inv} over the θ

$$M_{\min} \equiv \min_{\theta \in [0, \pi/2]} \min\{M_{\text{inv}}^+(\theta), M_{\text{inv}}^-(\theta)\}. \quad (\text{B.7})$$

Figure 12 shows the distributions of M_{\min} for the WZ and SUSY benchmark point event samples for 1000 parton-level events. The SUSY benchmark point **C350-100** involves the parameters:

$$m_{\tilde{\chi}_1^\pm} = m_{\tilde{\chi}_2^0} = 350 \text{ GeV}, \quad m_{\tilde{\chi}_1^0} = 100 \text{ GeV}. \quad (\text{B.8})$$

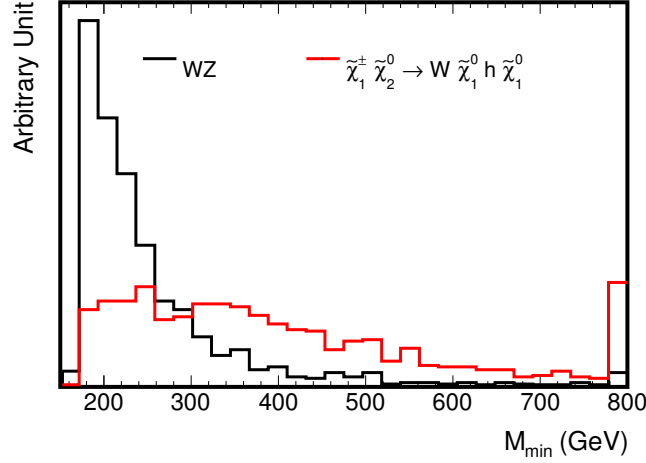


Figure 12. The M_{\min} distribution for WZ (black) and SUSY benchmark point (red) event samples.

C Calculation of the initial cross sections

For completeness we provide the branching ratios used to reproduce the initial cross sections that appear in table 4.

- **SUSY benchmark C350-100:** using Prospino 2.1, the NLO cross section for the SUSY benchmark point is $\sigma_{SUSY} \simeq 200 \text{ fb}$. For the signal, we consider the decays of the W to all three lepton families and the decays of the Higgs boson to either $\tau^+\tau^-$ or W^+W^- (again with the W s decaying to all leptons). We also assume that $BR(\tilde{\chi}_2^0 \rightarrow h\tilde{\chi}_1^0) = 1$. Hence:

$$\begin{aligned}
 \sigma(SUSY)_{\text{initial}} &= \sigma_{SUSY} \times BR(W \rightarrow \ell/\tau\nu) \\
 &\quad \times (BR(h \rightarrow \tau^+\tau^-) + BR(h \rightarrow W^+W^-) \times BR(W \rightarrow \ell/\tau\nu)^2) \\
 &= \sigma_{SUSY} \times 0.3257 \\
 &\quad \times (0.0632 + 0.2155 \times 0.3257^2) \\
 &\simeq 5.7 \text{ fb} .
 \end{aligned} \tag{C.1}$$

- **WZ :** we allow for $(W \rightarrow \ell\nu, Z \rightarrow \tau^+\tau^-)$ or $(W \rightarrow \tau\nu, Z \rightarrow \ell^+\ell^-)$. We use the NLO cross section $\sigma(WZ) = 51.82 \text{ pb}$, according to [67]. We obtain: $\sigma_{\text{initial}} = \sigma(WZ) \times (BR(W \rightarrow \ell\nu)BR(Z \rightarrow \tau^+\tau^-) + BR(W \rightarrow \tau\nu)BR(Z \rightarrow \ell^+\ell^-)) = (51.82 \times 10^3) \times (0.22 \times 3.37 \times 10^{-2} + 0.11 \times (6.7 \times 10^{-2})) \text{ pb} \simeq 767 \text{ fb}$.
- **W +jets:** the **AlpGen** tree-level cross section merged to the **HERWIG++** parton shower is $\sigma(W + \text{jets}) \simeq 300 \text{ pb}$ per lepton flavour (electrons, muons or taus). This was calculated for 2 associated partons with the W boson.
- **hV :** the processes hW and hZ are included. We assume $\sigma(hW) = 1504 \text{ fb}$ and $\sigma(hZ) = 883 \text{ fb}$ at NNLO QCD + NLO EW, taken from [68] for $M_h = 125 \text{ GeV}$.

We allow $h \rightarrow W^+W^-$ and $h \rightarrow \tau^+\tau^-$, impose no constrain on the W decays and allow for the Z to decay to all leptons. Hence: $\sigma_{\text{initial}} = (BR(h \rightarrow \tau^+\tau^-) + BR(h \rightarrow W^+W^-)) \times (\sigma(hZ) \times BR(Z \rightarrow \tau^+\tau^-/\ell^+\ell^-) + \sigma(hW)) = (0.0632 + 0.2155) \times (883 \times (10.1 \times 10^{-2}) + 1504) \text{ fb} \simeq 443 \text{ fb}$.

- $t\bar{t}h$, $t\bar{t}$: we assume that the efficiency of tagging jets originating from the decays b quarks is 70%. If one then vetoes events that contain at least one b -tagged jet, then for events containing $t\bar{t}$, a $1 - (0.7^2 + 2 \times 0.3 \times 0.7) = 0.09$ rejection factor can be achieved. We consider only leptonic decays of the W bosons originating from the decays of the top quarks and only consider $h \rightarrow \tau^+\tau^-$. We assume total cross sections: $\sigma(t\bar{t}) \sim 900 \text{ pb}$ and $\sigma(t\bar{t}h) \simeq 611 \text{ fb}$ [68]. This gives: $\sigma_{\text{initial}}(t\bar{t} \rightarrow \text{leptons} + \text{jets}) = 0.09 \times 900 \text{ pb} \times BR(W \rightarrow \ell/\tau\nu)^2 \simeq 8600 \text{ fb}$ and $\sigma_{\text{initial}}(t\bar{t}h \rightarrow \text{leptons} + \text{jets} + (\tau^+\tau^-)) \simeq 3.5 \text{ fb}$.
- Z +jets: the **AlpGen** tree-level cross section merged to the **HERWIG++** parton shower is $\sigma(Z + \text{jets}) \simeq 300 \text{ pb}$ per lepton flavour (electrons, muons or taus). This sample has been produced with one associated parton with the Z boson.

D Discovery with low statistics

Discovery occurs when the probability of obtaining a given experimental result, which contains some signal, is small when compared to the expected background hypothesis. How small this probability should be is somewhat a matter of preference and convention. Nowadays, in high energy physics, these probabilities are taken to correspond to 3 standard deviations away from the assumed central value of a Gaussian for the case of ‘evidence’ of a signal, and 5 standard deviations for the case of ‘discovery’ of a signal. On the other hand, exclusion is based on the probability of having fewer events than the background alone would give, given the signal plus background hypothesis.

To be concrete, let us assume that we are performing counting experiments of events, obtaining as a result, N_i counts in each experiment i . Let us assume that in one specific experiment, we obtained a measurement n_{obs} . By some theoretical prediction, for example obtained using a Monte Carlo event generator, or otherwise, the expected background number of events in this experiment is given to be b . We can assume that the counts N_i are random variables, distributed according to some distribution $P(N_i, b)$, where b is the mean of the distribution. In this case, the probability of obtaining n_{obs} or more events, when the mean is equal to the expected background b is given by:

$$P(N \geq n_{\text{obs}}, b) = \sum_{i=n_{\text{obs}}}^{i=\infty} P(N_i, b), \quad (\text{D.1})$$

where the sum can also be turned into an integral in the continuous variable case. In simple words, according to the ‘background only’ distribution, getting a measurement of n_{obs} or more amounts to the probability of the shaded area in figure 13., and this probability tells you how likely b is as an assumption of the mean of the distribution.

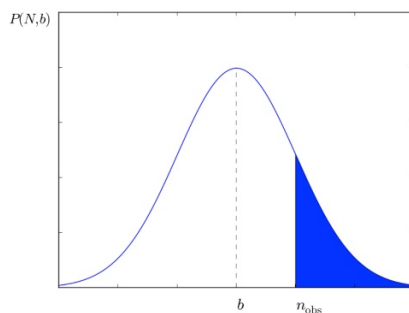


Figure 13. The shaded region in the above probability distribution shows the probability of obtaining $N > n_{\text{obs}}$ events.

In the specific case of the Poisson distribution:

$$\text{Pois}(N_i, b) = \frac{b^{N_i}}{N_i!} e^{-b}, \quad (\text{D.2})$$

then eq. (D.1) becomes $P(N \geq n_{\text{obs}}, b) = \sum_{i=n_{\text{obs}}}^{i=\infty} \frac{b^{N_i}}{N_i!} e^{-b}$. This sum can be shown (see appendix E) to be equivalent to the so-called ‘regularised incomplete gamma function’, $\Gamma_{\text{reg}}(s, x)$:

$$P(N \geq n_{\text{obs}}, b) = \sum_{i=n_{\text{obs}}}^{i=\infty} \frac{b^{N_i}}{N_i!} e^{-b} = \Gamma_{\text{reg}}(n_{\text{obs}}, b) = \Gamma(n_{\text{obs}}, b) / \Gamma(n_{\text{obs}}), \quad (\text{D.3})$$

for $n_{\text{obs}} > 0$, and where we have defined the ‘unregularised incomplete gamma function’:

$$\Gamma(n_{\text{obs}}, b) = \int_0^b dt \, t^{n_{\text{obs}}-1} e^{-t}, \quad (\text{D.4})$$

and $\Gamma(n_{\text{obs}})$ is defined in eq. (E.2) in the following section, for $n = n_{\text{obs}}$. We can then calculate the probability for discovery. This is given by $P(N \geq n_{\text{obs}}, b)$ for $n_{\text{obs}} = s + b$, where s is the expected signal contribution to the event counts. This probability will differ from the one obtained using the large sample (i.e. Gaussian) approximation, in which the significance is given approximately by $\sigma \sim s/\sqrt{b}$. For exclusion, we need to calculate the probability of having less than b events, under the assumption that the expected number of events is $s + b$, i.e. $P(N < b, s + b)$.

E Cumulative distribution for Poisson random variables

The unregularised incomplete gamma function is given by:

$$\Gamma(n, x) = \int_0^x dt \, t^{n-1} e^{-t}. \quad (\text{E.1})$$

One can then define the gamma function:

$$\Gamma(n) = \lim_{x \rightarrow \infty} \Gamma(n, x) = \int_0^\infty dt \, t^{n-1} e^{-t}. \quad (\text{E.2})$$

The integral that appears in $\Gamma(n, x)$ can be expanded by performing consecutive integrations by parts:

$$\begin{aligned}
\int_0^x dt \, t^{n-1} e^{-t} &= -e^{-t} t^{n-1} \Big|_0^x + (n-1) \int_0^x dt \, t^{n-2} e^{-t} \\
&= -e^{-x} x^{n-1} - (n-1) e^{-x} x^{n-2} \\
&\quad + (n-1)(n-2) \int_0^x dt \, t^{n-3} e^{-t} , \\
&= -e^{-x} [x^{n-1} + (n-1)x^{n-2} \\
&\quad + (n-1)(n-2)x^{n-3} + \dots] \\
&\quad + (n-1)(n-2)(n-3) \dots (1) [1 - e^{-x}].
\end{aligned}
\tag{E.3}$$

From the last equality in the above expression we can deduce that

$$\Gamma(n) = (n-1)! . \tag{E.4}$$

For $n > 0$, dividing $\Gamma(n, x)$ by $\Gamma(n)$, we obtain:

$$\begin{aligned}
\frac{\Gamma(n, x)}{\Gamma(n)} &= 1 - e^{-x} \left[\frac{x^{n-1}}{(n-1)!} + \frac{x^{n-2}}{(n-2)!} + \frac{x^{n-3}}{(n-3)!} + \dots \right] \\
&= 1 - \sum_{i \leq n-1} \frac{x^i e^{-x}}{i!} = \sum_{i=n}^{\infty} \frac{x^i e^{-x}}{i!} ,
\end{aligned}
\tag{E.5}$$

which is nothing but the cumulative sum for the Poisson distribution.

F Individual signal regions

In figures 14 and 15 we demonstrate the individual signal regions contributing to the envelopes shown in figures 8 and 9. The analyses at each luminosity are identical and the more ‘irregular’ form at lower luminosities is related to the Poisson statistics that govern the smaller number of events in those cases.

The sensitivity of the BSM searches depends on the hardness of the final state particles but not strongly on the angular distributions and polarisation. The hardness of the final state particles is primarily determined by the topology and mass spectrum of the BSM particles appearing in the process. Therefore, for example, if the LSP neutralino is replaced by a gravitino with the same mass, keeping the rest of the topology unchanged, the signal region efficiency is almost unchanged [69].

In table 6 we show the resulting cross sections after applying each of the signal regions, defined in table 5. These can be used in conjunction with the efficiency data files for the signal on the M_2 - M_1 plane attached to this article¹³ to construct the signal cross section for each signal region for explicit BSM scenarios with $\tilde{N}\tilde{C}^\pm \rightarrow (h\chi)(W^\pm\chi)$ topology, where \tilde{N}

¹³The file corresponding to signal region X is “efficiency_regionX.expanded.dat”, located in the subdirectory “efficiencies” of the distribution of this article.

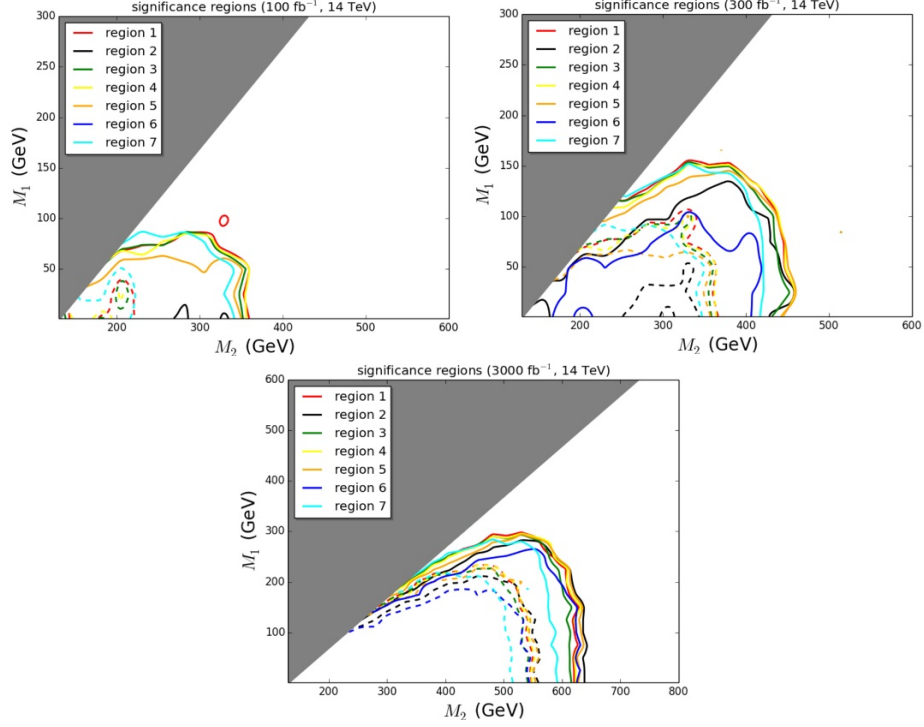


Figure 14. The significance on the M_2 - M_1 plane obtained for the signal regions defined in table 5 at integrated luminosities of 100 fb^{-1} (upper left), 300 fb^{-1} (upper right) and 3000 fb^{-1} (bottom). The solid curves show the 3σ evidence region, whereas the dashed curves show the 5σ discovery region.

signal region	SR1	SR2	SR3	SR4	SR5	SR6	SR7
$\sigma_{\text{bkg}} \text{ (fb)}$	0.312	0.189	0.338	0.297	0.254	0.195	0.512

Table 6. The resulting sum of cross sections for the backgrounds for the different signal regions (SR) used in the analysis.

and \tilde{C}^\pm are massive BSM particles with the same mass, M_2 , and χ is an invisible particle with mass M_1 . This feature makes simplified models very useful for recasting analysis. One can calculate the signal cross section for the process in question according to the given model:

$$[\text{signal efficiency, signal region X}] \times [\text{signal cross section}] \times [\text{BR}] , \quad (\text{F.1})$$

and use this in conjunction with the background cross section for region X as given in the table to obtain the p-value over the parameter space. Our efficiency data considers only the process with the $W \rightarrow \ell/\tau, \nu$ and Higgs bosons decaying inclusively to leptons (either $h \rightarrow \tau^+\tau^-$ or $h \rightarrow W^+W^- \rightarrow (e^+e^-, \mu^+\mu^-, \tau^+\tau^-) + \cancel{E}_T$). The [BR] factor in eq. (F.1) should therefore include these branching ratios.

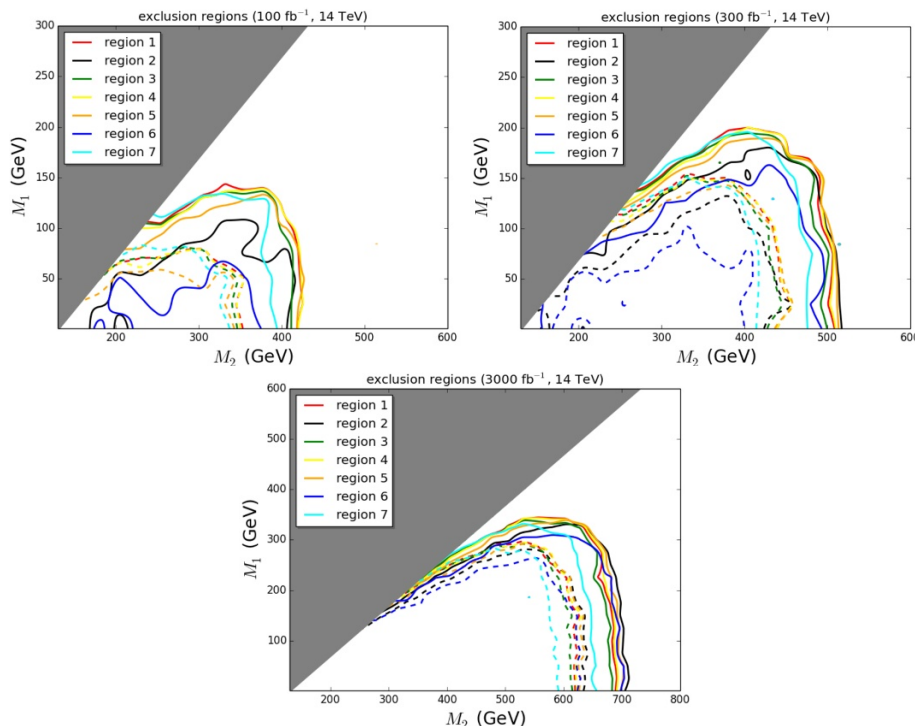


Figure 15. The exclusion on the M_2 - M_1 plane obtained for the signal regions defined in table 5 at integrated luminosities of 100 fb^{-1} (upper left), 300 fb^{-1} (upper right) and 3000 fb^{-1} (bottom). The solid curves show the 2σ exclusion boundary, whereas the dashed curves show the 3σ boundary.

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